

「電磁波工学の基礎」章末問題解答

* 解答が自明なものは省略します.

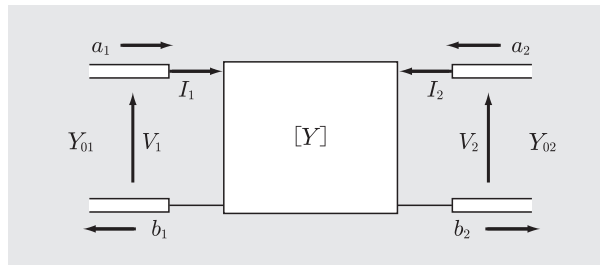
■ 1章の問題

□ 1~4 省略

■ 2章の問題

□ 1 省略

□ 2



$$\begin{pmatrix} I_1 \\ I_2 \end{pmatrix} = \begin{pmatrix} Y_{11} & Y_{12} \\ Y_{21} & Y_{22} \end{pmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix} \quad (1)$$

式(2.86)より

$$\begin{pmatrix} V_1 \\ V_2 \end{pmatrix} = \begin{pmatrix} \frac{1}{\sqrt{Y_{01}}}(a_1 + b_1) \\ \frac{1}{\sqrt{Y_{02}}}(a_2 + b_2) \end{pmatrix}, \quad (2)$$

$$\begin{pmatrix} I_1 \\ I_2 \end{pmatrix} = \begin{pmatrix} \sqrt{Y_{01}}(a_1 - b_1) \\ \sqrt{Y_{02}}(a_2 - b_2) \end{pmatrix}$$

式(2)を式(1)へ代入して

$$\begin{pmatrix} \sqrt{Y_{01}}(a_1 - b_1) \\ \sqrt{Y_{02}}(a_2 - b_2) \end{pmatrix} = \begin{pmatrix} Y_{11} & Y_{12} \\ Y_{21} & Y_{22} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{Y_{01}}}(a_1 + b_1) \\ \frac{1}{\sqrt{Y_{02}}}(a_2 + b_2) \end{pmatrix}, \quad (3)$$

$$\begin{pmatrix} \sqrt{Y_{01}} & 0 \\ 0 & \sqrt{Y_{02}} \end{pmatrix} \begin{pmatrix} a_1 - b_1 \\ a_2 - b_2 \end{pmatrix} = (Y) \begin{pmatrix} \frac{1}{\sqrt{Y_{01}}} & 0 \\ 0 & \frac{1}{\sqrt{Y_{02}}} \end{pmatrix} \begin{pmatrix} a_1 + b_1 \\ a_2 + b_2 \end{pmatrix}, \quad (4)$$

$$\begin{pmatrix} \sqrt{Y_{01}} & 0 \\ 0 & \sqrt{Y_{02}} \end{pmatrix}^{-1} = \frac{1}{\sqrt{Y_{01}Y_{02}}} \begin{pmatrix} \sqrt{Y_{02}} & 0 \\ 0 & \sqrt{Y_{01}} \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{\sqrt{Y_{01}}} & 0 \\ 0 & \frac{1}{\sqrt{Y_{02}}} \end{pmatrix} = \begin{pmatrix} \sqrt{Z_{01}} & 0 \\ 0 & \sqrt{Z_{02}} \end{pmatrix}, \quad (5)$$

だから、式(5)を式(4)の左からかけて

$$[a] - [b] = [\sqrt{Z_0}] [Y] [\sqrt{Z_0}] \{[a] + [b]\},$$

$$[\hat{Y}] \triangleq [\sqrt{Z_0}] [Y] [\sqrt{Z_0}]$$

とすれば,

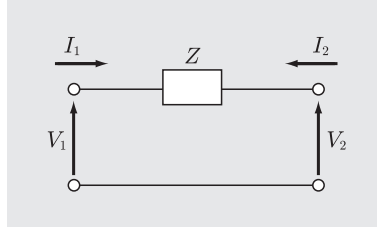
$$[a] - [b] = [\hat{Y}] \{[a] + [b]\},$$

$$\{[1] - [\hat{Y}]\}[a] = \{[1] + [\hat{Y}]\}[b]$$

ゆえに

$$[b] = \underbrace{\{[1] + [\hat{Y}]\}^{-1} \{[1] - [\hat{Y}]\}}_{[S]} [a]$$

□ 3



$$\begin{cases} I_1 = -I_2 \\ V_1 = I_1 Z + V_2 \end{cases} \implies \begin{cases} I_1 = \frac{I}{Z}(V_1 - V_2) \\ I_2 = -\frac{I}{Z}(V_1 - V_2) \end{cases}$$

ゆえに

$$\begin{pmatrix} I_1 \\ I_2 \end{pmatrix} = \frac{1}{Z} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix},$$

$$[\hat{Y}] = [\sqrt{Z_0}] [Y] [\sqrt{Z_0}]$$

$$= \begin{pmatrix} \sqrt{Z_0} & 0 \\ 0 & \sqrt{Z_0} \end{pmatrix} \frac{1}{Z} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} \sqrt{Z_0} & 0 \\ 0 & \sqrt{Z_0} \end{pmatrix}$$

$$= \frac{Z_0}{Z} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \frac{1}{Z} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix}$$

よって,

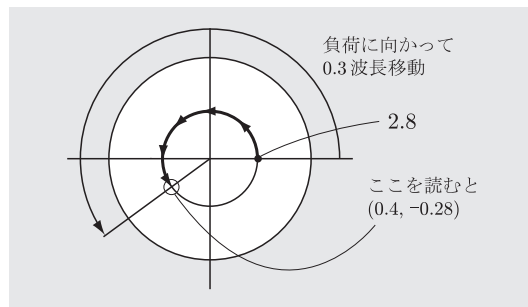
$$\begin{aligned}
 [S] &= \left\{ \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{\hat{Z}} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \right\}^{-1} \left\{ \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \frac{1}{\hat{Z}} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \right\} \\
 &= \begin{pmatrix} 1 + \frac{1}{\hat{Z}} & -\frac{1}{\hat{Z}} \\ -\frac{1}{\hat{Z}} & 1 + \frac{1}{\hat{Z}} \end{pmatrix}^{-1} \begin{pmatrix} 1 - \frac{1}{\hat{Z}} & \frac{1}{\hat{Z}} \\ \frac{1}{\hat{Z}} & 1 - \frac{1}{\hat{Z}} \end{pmatrix} \\
 &= \frac{1}{\left(1 + \frac{1}{\hat{Z}}\right)^2 - \left(\frac{1}{\hat{Z}}\right)^2} \begin{pmatrix} 1 + \frac{1}{\hat{Z}} & \frac{1}{\hat{Z}} \\ \frac{1}{\hat{Z}} & 1 + \frac{1}{\hat{Z}} \end{pmatrix}^{-1} \begin{pmatrix} 1 - \frac{1}{\hat{Z}} & \frac{1}{\hat{Z}} \\ \frac{1}{\hat{Z}} & 1 - \frac{1}{\hat{Z}} \end{pmatrix} \\
 &= \frac{1}{1 + \frac{1}{\hat{Z}}} \begin{pmatrix} 1 & \frac{2}{\hat{Z}} \\ \frac{2}{\hat{Z}} & 1 \end{pmatrix} \\
 &= \frac{1}{\hat{Z} + 2} \begin{pmatrix} \hat{Z} & 2 \\ 2 & \hat{Z} \end{pmatrix}
 \end{aligned}$$

□ 4 (1) 線路上の波長は

$$\lambda_g = \frac{3 \times 10^8}{2 \times 10^9} \times \frac{1}{\sqrt{6.25}} = 60 \text{ [mm]}$$

従って 18 mm は

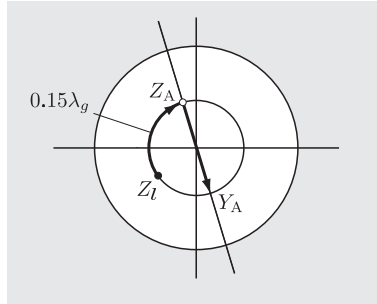
$$\frac{18}{60} = 0.3 \text{ [波長]}$$



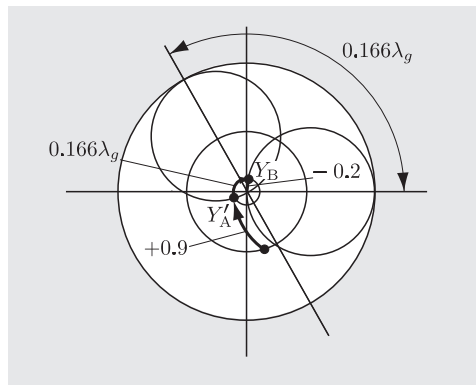
ゆえに、アンテナインピーダンスは

$$75 \times (0.4 - j0.28) = 30 - j21 \text{ [\Omega]}$$

(2) 9 mm \rightarrow $0.15\lambda_g$, 10 mm \rightarrow $0.166\lambda_g$. スミスチャート上で $Z_l \rightarrow Z_A \rightarrow Y_A$ と進む.



Y_A に $+0.9$ のサセプタンスをつなぎ、 Y'_A へ移動。そこから $0.166\lambda_g$ (10 mm) さらにさかのぼり、 $G = 1$ の円上に到達できる (Y_B)。そこで、 -0.2 のサセプタンスをつなぐと、原点へ到達する。

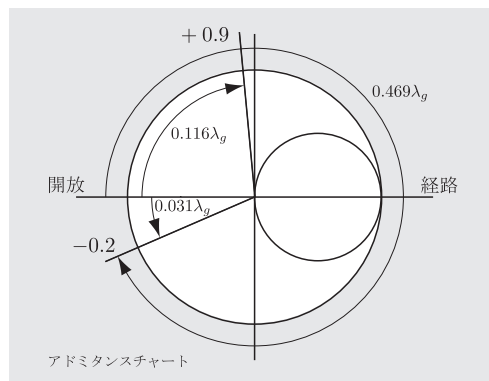


ゆえに

A 点に 0.469×60 [mm] = 28 [mm] (-0.2 の正規化サセプタンスに相当)

B 点に 0.116×60 [mm] = 7 [mm] ($+0.9$ の正規化サセプタンスに相当)

の終端開放線路をつなげばよい。



■ 3章の問題

□ 1 式(3.44)の導出.

$$\begin{aligned}
 \mathbf{S} &\triangleq \mathbf{E} \times \mathbf{H}^* \quad (\text{複素ポインティングベクトル}) \\
 \nabla \cdot \mathbf{S} &= \nabla \cdot (\mathbf{E} \times \mathbf{H}^*) = \mathbf{H}^* \cdot \nabla \times \mathbf{E} - \mathbf{E} \cdot \nabla \times \mathbf{H}^* \\
 &\stackrel{(3.14),(3.15)}{=} \mathbf{H}^* \cdot (-j\omega\mu\mathbf{H}) - \mathbf{E} \cdot (\sigma - j\omega\varepsilon)\mathbf{E}^* \\
 &= -j\omega\mu|\mathbf{H}|^2 - \sigma|\mathbf{E}|^2 + j\omega\varepsilon|\mathbf{E}|^2
 \end{aligned}$$

ゆえに

$$\int_V \nabla \cdot \mathbf{S} dV = - \int_V \sigma |\mathbf{E}|^2 dV + j\omega \int_V (\varepsilon |\mathbf{E}|^2 - \mu |\mathbf{H}|^2) dV$$

一方,

$$\int_V \nabla \cdot \mathbf{S} dV \stackrel{\text{ガウスの定理}}{=} \int_S \mathbf{S} d\mathbf{S} = \int_S (\mathbf{E} \times \mathbf{H}^*) d\mathbf{S}$$

よって

$$\int_S (\mathbf{E} \times \mathbf{H}^*) d\mathbf{S} + \int_V \sigma |\mathbf{E}|^2 dV = j\omega \int_V (\varepsilon |\mathbf{E}|^2 - \mu |\mathbf{H}|^2) dV$$

□ 2 省略

□ 3

● 遮断周波数

$$f_c = \frac{\omega_c}{2\pi} = \frac{1}{2\pi} ck_c = \frac{c}{2\pi} \sqrt{\left(\frac{\pi}{a}\right)^2 + \left(\frac{\pi}{b}\right)^2} \quad (\text{TM}_{11} \rightarrow m = n = 1)$$

● 位相定数

$$\beta = \sqrt{\omega^2 \varepsilon \mu - k_c^2} = \sqrt{\frac{\omega^2}{c^2} - \frac{\pi^2}{a^2} - \frac{\pi^2}{b^2}} \quad (\text{または} = \pi \sqrt{\left(\frac{2f}{c}\right)^2 - \frac{1}{a^2} - \frac{1}{b^2}})$$

● 管内波長

$$\lambda_g = \frac{2\pi}{\beta} = \frac{2}{\sqrt{\left(\frac{2f}{c}\right)^2 - \frac{1}{a^2} - \frac{1}{b^2}}} \quad (\text{または} = \frac{2\pi}{\sqrt{\frac{\omega^2}{c^2} - \frac{\pi^2}{a^2} - \frac{\pi^2}{b^2}}})$$

● 位相速度

$$v_p = \frac{\omega}{\beta} = \frac{\omega}{\sqrt{\frac{\omega^2}{c^2} - \frac{\pi^2}{a^2} - \frac{\pi^2}{b^2}}} \quad (\text{または} = \frac{1}{\sqrt{\frac{1}{c^2} - \frac{1}{\omega^2} \left(\frac{\pi^2}{a^2} + \frac{\pi^2}{b^2}\right)}})$$

● 群速度

$$\begin{aligned}
 v_g &= \left(\frac{d\beta}{d\omega}\right)^{-1} = \left(\frac{1}{2}(\omega^2 \varepsilon \mu - k_c^2)^{-\frac{1}{2}} 2\omega \varepsilon \mu\right)^{-1} \\
 &= \frac{\sqrt{\omega^2 \varepsilon \mu - k_c^2}}{\omega \varepsilon \mu} = \frac{\sqrt{\frac{\omega^2}{c^2} - \frac{\pi^2}{a^2} - \frac{\pi^2}{b^2}}}{\omega \varepsilon \mu} = \frac{c^2}{\omega} \sqrt{\frac{\omega^2}{c^2} - \frac{\pi^2}{a^2} - \frac{\pi^2}{b^2}} \\
 &\quad (\text{または} = \frac{c^2}{2f} \sqrt{\frac{4f^2}{c^2} - \frac{1}{a^2} - \frac{1}{b^2}})
 \end{aligned}$$

• 伝送電力

$$\begin{aligned}
\mathbf{k} \cdot S_z &= \mathbf{E}_t \times \mathbf{H}_t^* = (\mathbf{i}E_x + \mathbf{j}E_y) \times (\mathbf{i}H_x^* + \mathbf{j}H_y^*) = \mathbf{k}(E_x H_y^* - E_y H_x^*) \\
&= \mathbf{k} \left(-\frac{\gamma k_x}{k_c^2} E_{11} \cos k_x x \sin k_y y \cdot \frac{j\omega\epsilon k_x}{k_c^2} E_{11}^* \cos k_x x \sin k_y y \right. \\
&\quad \left. - (-) \frac{\gamma k_y}{k_c^2} E_{11} \sin k_x x \cos k_y y \cdot (-) \frac{j\omega\epsilon k_y}{k_c^2} E_{11}^* \sin k_x x \cos k_y y \right) \\
&= \mathbf{k} \left(\frac{\beta\omega\epsilon k_x^2}{k_c^4} |E_{11}|^2 \cos^2 k_x x \sin^2 k_y y + \frac{\beta\omega\epsilon k_y^2}{k_c^4} |E_{11}|^2 \sin^2 k_x x \cos^2 k_y y \right) \\
&= \mathbf{k} \left(\frac{\beta\omega\epsilon |E_{11}|^2}{k_c^4} (k_x^2 \cos^2 k_x x \sin^2 k_y y + k_y^2 \sin^2 k_x x \cos^2 k_y y) \right) \\
&= \mathbf{k} \frac{\beta\omega\epsilon |E_{11}|^2}{k_c^4} \left(\left(\frac{\pi}{a}\right)^2 \cos^2 \frac{\pi x}{a} \sin^2 \frac{\pi y}{b} + \left(\frac{\pi}{b}\right)^2 \sin^2 \frac{\pi x}{a} \cos^2 \frac{\pi y}{b} \right)
\end{aligned}$$

$$\begin{aligned}
\therefore P &= \int_0^a \int_0^b \operatorname{Re} S_z dx dy \\
&= \frac{\beta\omega\epsilon |E_{11}|^2}{k_c^4} \left[\left(\frac{\pi}{a}\right)^2 \int_0^a \int_0^b \cos^2 \frac{\pi x}{a} \sin^2 \frac{\pi y}{b} dx dy \right. \\
&\quad \left. + \left(\frac{\pi}{b}\right)^2 \int_0^a \int_0^b \sin^2 \frac{\pi x}{a} \cos^2 \frac{\pi y}{b} dx dy \right] \\
&= \frac{\beta\omega\epsilon |E_{11}|^2}{\left(\left(\frac{\pi}{a}\right)^2 + \left(\frac{\pi}{b}\right)^2\right)^2} \int_0^a \int_0^b \left\{ \left(\frac{\pi}{a}\right)^2 \cos^2 \frac{\pi x}{a} \sin^2 \frac{\pi y}{b} \right. \\
&\quad \left. + \left(\frac{\pi}{b}\right)^2 \sin^2 \frac{\pi x}{a} \cos^2 \frac{\pi y}{b} \right\} dx dy
\end{aligned}$$

• 特性インピーダンス

$$\text{TM} : Z_E = \frac{\beta}{\omega\epsilon} = \frac{1}{\omega\epsilon} \sqrt{\omega^2 - \frac{\pi^2}{a^2} - \frac{\pi^2}{b^2}}$$

□ 4 (1) Cバンドは4~8 GHz. 矩形導波管の辺を a, b として, 遮断周波数は式 (3.82) より $f_c = \frac{c}{2} \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2}$.

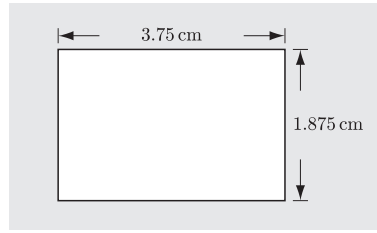
$n \backslash m$	0	1	2
0	-	$\frac{c}{2b}$	$\frac{2c}{2b}$
1	$\frac{c}{2a}$	$\frac{c}{2} \sqrt{\left(\frac{1}{a}\right)^2 + \left(\frac{1}{b}\right)^2}$	$\frac{c}{2} \sqrt{\left(\frac{1}{a}\right)^2 + \left(\frac{2}{b}\right)^2}$
2	$\frac{2c}{2a}$	$\frac{c}{2} \sqrt{\left(\frac{2}{a}\right)^2 + \left(\frac{1}{b}\right)^2}$	$\frac{c}{2} \sqrt{\left(\frac{2}{a}\right)^2 + \left(\frac{2}{b}\right)^2}$

ゆえに、4~8 GHz で TE₁₀ の単一モード導波管とするには

$$\frac{c}{2a} = 4 \times 10^9, \quad \frac{c}{2b} = 8 \times 10^9$$

とすればよい.

$$\therefore a = \frac{3 \times 10^8}{2 \times 4 \times 10^9} = 3.75 \text{ [cm]}, \quad \therefore b = \frac{3 \times 10^8}{2 \times 8 \times 10^9} = 1.875 \text{ [cm]}$$



(2) $\frac{4+8}{2} = 6$ GHz での管内波長 λ_g , 特性インピーダンス Z_0 , 群速度 v_g を求める.

$$\lambda = \frac{3 \times 10^8}{6 \times 10^9} = 5 \text{ [cm]}, \quad \lambda_c = \frac{3 \times 10^8}{4 \times 10^9} = 7.5 \text{ [cm]}$$

式 (3.96) より

$$\lambda_g = \frac{\lambda}{\sqrt{1 - \frac{\lambda}{\lambda_c}}} = \frac{5}{\sqrt{1 - \left(\frac{5}{7.5}\right)^2}} = 6.71 \text{ [cm]}$$

式 (3.104) より

$$\begin{aligned} Z_0 &= \frac{j\omega\mu}{\gamma} = \frac{j\omega\mu}{j\beta} = \frac{\omega\mu}{\beta} = \frac{2\pi f\mu}{\frac{2\pi}{\lambda_g}} = \lambda_g \cdot f \cdot \mu \\ &= 6.71 \times 10^{-2} \times 6 \times 10^9 \times \underbrace{4\pi \cdot 10^{-7}}_{\text{式 (3.33)}} = 505.8 \text{ [\Omega]} \end{aligned}$$

$$v_g = \frac{c^2}{v_p} = \frac{c^2}{f \cdot \lambda_g} = \frac{(3 \times 10^8)^2}{(6 \times 10^9) \cdot (6.71 \times 10^{-2})} = 2.24 \times 10^8 \text{ [m/s]}$$

(3) Xバンドは8~12 GHz. $f_c = \frac{c}{2} \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2}$ に $a = 3.75$ [cm], $b = 1.875$ [cm] を代入して, 下表を得る.

$m \backslash n$	0	1	2	3
0	-	8**	16	24
1	4*	8.94**	16.49	24.33
2	8**	11.3**	17.9	25.3
3	12	14.4	20	26.8

従って、TE₁₀ (*) に加え、TE₀₁, TE₁₁, TM₁₁, TE₂₀, TE₂₁, TM₂₁ (**) が励振され得る。従って、Xバンドでは7つのモードが存在し得る。

□ 5 (1)

$$Z_0 = 138 \log_{10} \frac{b}{a} = 138 \log_{10} \frac{\frac{d}{2}}{\frac{d}{2}} = 138 \log_{10} \frac{d}{2} \implies 100 [\Omega]$$

$$\therefore \frac{100}{138} = \log_{10} \frac{d}{2}, \quad \therefore 10^{\frac{100}{138}} = \frac{d}{2}, \quad \therefore d = 2 \times 10^{\frac{100}{138}} = 10.6 [\text{mm}]$$

(2)

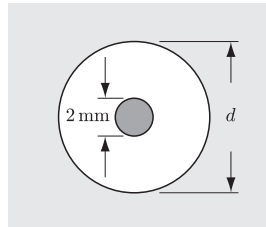
$$f_c \cong \frac{c}{\pi} \cdot \frac{1}{a+b} = \frac{c}{\pi} \cdot \frac{1}{\left(\frac{2}{2} + \frac{10.6}{2}\right) \times 10^{-3}} = 15.2 [\text{GHz}]$$

(3)

$$E_r|_a = \frac{V_0}{a} = 2 \times 10^5 [\text{V/m}], \quad \therefore V_0 = 2 \times 10^5 \times \underbrace{\frac{2}{2}}_a \times 10^{-3} = 2 \times 10^2$$

$$\therefore P = \frac{2\pi}{\zeta} |V_0|^2 \ln \frac{b}{a} = \frac{2\pi}{377} |200|^2 \ln \frac{10.6}{2} = 1.1 \times 10^3 [\text{W}] = 1.1 [\text{kW}]$$

(平均ではこの半分の 555 W。どちらも正解)

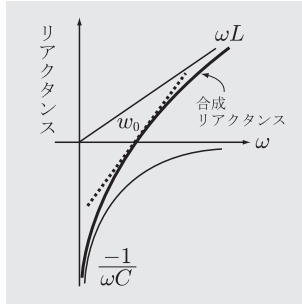


□ 6 図 3.34, 図 3.33 より $\frac{a}{b} = 3$ の Z_0 と $\sqrt{\epsilon_{\text{eff}}}$ を求めると、各々、 70Ω と 2.5. よって特性インピーダンスは $\frac{70}{2.5} = 30 [\Omega]$.

■ 4章の問題

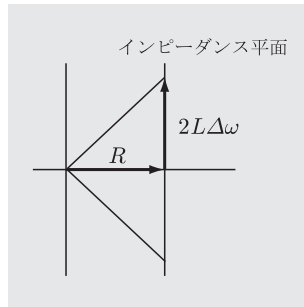
□ 1

$$\text{共振: } \omega_0 L - \frac{1}{\omega_0 C} = 0 \implies \omega_0 = \frac{1}{\sqrt{LC}}$$



上図の破線で示した直線（接線）は

$$\frac{d}{d\omega} \left(\omega L - \frac{1}{\omega C} \right) \Big|_{\omega=\omega_0} \cdot \frac{\Delta\omega}{(\omega-\omega_0)} = \left(L + \frac{1}{C}\omega^{-2} \right) \Big|_{\omega=\omega_0} \cdot \Delta\omega = 2L\Delta\omega$$



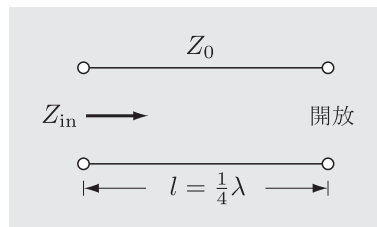
図より

$$2L\Delta\omega = R, \quad \therefore \Delta\omega_0 = \frac{R}{2L}$$

よって

$$Q \triangleq \frac{\omega_0}{2\Delta\omega_0} = \frac{1}{\sqrt{LC}} \cdot \frac{1}{2} \cdot \frac{2L}{R} = \frac{1}{R} \sqrt{\frac{L}{C}}$$

□ 2



式(2.73)より

$$Z_{in} = -jZ_0 \cot \beta l$$

一般に, $\coth Z = \frac{1}{\gamma} \cot \frac{Z}{j}$ であることから

$$Z_{\text{in}} = Z_0 \coth \gamma l$$

これに, $l = \frac{\lambda_{g0}}{4} = \frac{1}{4} \cdot \frac{2\pi C}{\omega_0}$ を代入すると $\beta l = \frac{\pi}{2} \left(1 + \frac{\Delta\omega}{\omega_0}\right)$ なので

$$\begin{aligned} Z_{\text{in}} &= Z_0 \coth \gamma l = Z_0 \coth \left(\alpha l + j \frac{\pi}{2} \left(1 + \frac{\Delta\omega}{\omega_0}\right) \right) \\ &= Z_0 \coth \left(\alpha l + j \frac{\pi}{2} \cdot \frac{\Delta\omega}{\omega_0} + j \frac{\pi}{2} \right) \\ &= Z_0 \tanh \left(\alpha l + j \frac{\pi}{2} \cdot \frac{\Delta\omega}{\omega_0} \right) \\ &\cong Z_0 \operatorname{sech}^2(0) \cdot \left(\alpha l + j \frac{\pi}{2} \cdot \frac{\Delta\omega}{\omega_0} \right) \quad (\text{if } \alpha \text{ 小} : \Delta\omega \rightarrow 0) \\ &= Z_0 \left(\alpha l + j \frac{\pi}{2} \cdot \frac{\Delta\omega}{\omega_0} \right) \\ &= \frac{\pi\alpha}{2\beta_0} Z_0 \left(1 + j \frac{\beta_0}{2\alpha} \left(\frac{2\Delta\omega}{\omega_0} \right) \right) \quad (l = \frac{\pi}{2\beta_0} \text{ を代入}) \end{aligned}$$

これを式(4.2)と比較すると ($\frac{\beta_0}{2\alpha} \triangleq Q$ として)

$$Z_{\text{in}} = \frac{\pi Z_0}{4Q} \left(1 + jQ \frac{2\Delta\omega}{\omega_0} \right)$$

を得る. 等価回路の対応は

$$\begin{cases} R = \frac{\pi Z_0}{4Q} \\ L = \frac{QR}{\omega_0} = \frac{\pi Z_0}{4\omega_0} \\ C = \frac{1}{\omega_0^2 L} = \frac{4\omega_0}{\pi Z_0} \cdot \frac{1}{\omega_0^2} = \frac{4}{\omega_0 \pi Z_0} \end{cases}$$

となる.

□ **3~6** 省略

□ **7** TE_{10n} について, 式(4.61)より, 共振波長は $\frac{1}{\left(\frac{\lambda}{2}\right)^2} = \frac{1}{a^2} + \frac{1}{\left(\frac{c}{n}\right)^2}$ で決まる. 共振周波数に直すと

$$f = \frac{\text{光速}}{\lambda} = \frac{\text{光速}}{2} \sqrt{\frac{1}{a^2} + \frac{1}{\left(\frac{c}{n}\right)^2}} \quad (4.51)$$

今, $f = 1.2f_c$ なので

$$\begin{aligned} \frac{\text{光速}}{2} \sqrt{\frac{1}{a^2} + \frac{1}{\left(\frac{c}{n}\right)^2}} &= \frac{\text{光速}}{2a} \cdot 1.2 \\ \therefore \frac{(1.2)^2}{a^2} &= \frac{1}{a^2} + \frac{1}{\left(\frac{c}{n}\right)^2}, \quad \therefore \left(\frac{c}{n}\right)^2 = \frac{a^2}{(1.2)^2 - 1} \end{aligned}$$

ゆえに

$$\frac{c}{n} = \frac{a}{\sqrt{(1.2)^2 - 1}} = \frac{a}{0.663} = \frac{5.7}{0.663} = 8.59 \text{ [mm]}$$

よって, TE_{101} について $c = 8.59 \text{ [mm]}$ とすればよい.

□ 8 (1) 図 3.34 より $\frac{a}{b} = 2$ で $\epsilon_r = 1$ の Z_0 を読むと 90Ω . 一方, 図 3.33 で, $\frac{a}{b} = 2$ で $\epsilon_r = 4$ の $\sqrt{\epsilon_{\text{eff}}}$ を読むと 1.8. ゆえに

$$\sqrt{\epsilon_{\text{eff}}} = \underline{3.24}$$

よって

$$\frac{90}{1.8} = \underline{50 \text{ [\Omega]}}$$

(2) 式 (4.28) より

$$\frac{\lambda_g}{2} - 2v_p Z_0 C = l$$

ゆえに

$$\begin{aligned} \lambda_g &= 2(l + 2v_p Z_0 C) = 2(7.5 \times 10^{-3} + 2 \times \frac{3 \times 10^8}{1.8} \times 50 \times 0.05 \times 10^{-12}) \\ &= 16.7 \text{ [mm]} \end{aligned}$$

ゆえに

$$f = \frac{C}{\sqrt{\epsilon_{\text{eff}}} \lambda_g} = \frac{3 \times 10^8}{1.8 \times 16.7 \times 10^{-3}} = \underline{10 \text{ [GHz]}}$$

(3)(4) 下図参照.

